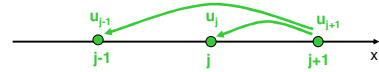


Numerical approximations of derivatives and integrals

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An implicit differencing scheme with second order accuracy



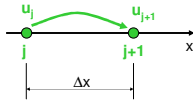
$$u_j = u_{j+1} + u'_{j+1}(-\Delta x) + u''_{j+1} \frac{\Delta x^2}{2} + o(\Delta x^2)$$

$$u_{j-1} = u_{j+1} + u'_{j+1}(-2\Delta x) + u''_{j+1} 2\Delta x^2 + o(\Delta x^2)$$

$$u_j - \frac{u_{j-1}}{4} = \frac{3}{4}u_{j+1} + u'_{j+1} \left(-\frac{\Delta x}{2}\right) + o(\Delta x^2)$$

$$u'_{j+1} = \frac{\frac{3}{2}u_{j+1} - 2u_j + \frac{1}{2}u_{j-1}}{\Delta x} + o(\Delta x)$$

Euler method



Taylor polynomial of the solution from point j to point j+1:

$$u_{j+1} = u_j + u'_j \Delta x + o(\Delta x)$$

This is an integration scheme of first order accuracy.

A differencing scheme with first order accuracy:

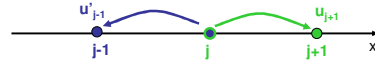
$$u'_j = \frac{u_{j+1} - u_j}{\Delta x} + o(1)$$

Another first order scheme:

$$u_j = u_{j+1} + u'_{j+1}(-\Delta x) + o(\Delta x)$$

u'_{j+1} usually is a given (but more complicated) function of x_{j+1} and u_{j+1} . Substitution of this function into the above formula leads to a more complicated expression for u_{j+1} . This kind of scheme is called **implicit**.

Adams-Basforth scheme



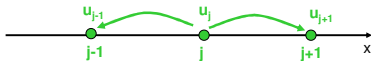
$$u_{j+1} = u_j + u'_j \Delta x + u''_j \frac{\Delta x^2}{2} + o(\Delta x^2)$$

$$u'_{j-1} = u'_j + u''_j(-\Delta x) + o(\Delta x) \quad \Bigg/ \quad + \dots \times \frac{\Delta x}{2}$$

$$u_{j+1} = u_j + \frac{3}{2}u'_j \Delta x - \frac{1}{2}u'_{j-1} \Delta x + o(\Delta x^2)$$

An explicit integrating scheme with second order accuracy. It is often used for **integrating the Navier-Stokes equations**.

CDS



$$u_{j+1} = u_j + u'_j \Delta x + u''_j \frac{\Delta x^2}{2} + o(\Delta x^2)$$

$$u_{j-1} = u_j + u'_j(-\Delta x) + u''_j \frac{\Delta x^2}{2} + o(\Delta x^2)$$

$$u'_j = \frac{u_{j+1} - u_{j-1}}{2\Delta x} + o(\Delta x)$$

Spatial derivatives in finite volume methods

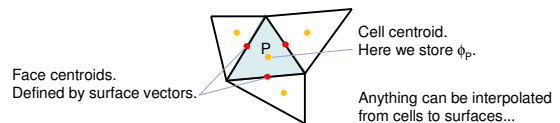
The generic transport equation in integral form:

$$\frac{\partial \rho \phi}{\partial t} + \nabla \cdot (\rho \phi \vec{v}) = \nabla \cdot \vec{S}_A + \nabla \cdot (\Gamma \nabla \phi) + S_v$$

In which ϕ is the mass concentration of a conserved quantity (eg. in kg/kg).

Spatial derivatives are always in $\text{div}(\dots)$, $\text{grad}(\dots)$ or $\text{div}(\text{grad}(\dots))$ forms. We only need to look for the discrete approximations of these operators, which is done - in the case of finite volume method - on the basis of surface and volume integrals along with some spatial interpolations.

The numerical mesh around the cell having its center in point P:



Approximation of the divergence operator

From the volume integral the divergence operator we can obtain an average value for the numerical cell.

The Gauss-Ostrogradskij theorem for a vector quantity \underline{u} :

$$\int_V \nabla \cdot \underline{u} dV = \oint_A \underline{u} \cdot d\underline{A}$$

For simplicity, we denote components of \underline{u} vector by u_i . The cell-average of the divergence operator is now:

$$\tilde{\nabla} \cdot \underline{u}_i = \frac{\sum_k \int_{A_k} u_{\perp} dA}{V_p}$$

in which ahol A_k a cella oldalfalainak indexe. The surface integral for one face is a scalar product:

$$\int_{A_k} u_{\perp} dA = \sum_{i=1}^3 u_i dA_i \quad \text{in which } u_i \text{ is one component of } \underline{u} \text{ interpolated to the cell surface.}$$

Gradient

A direct consequence of the Gauss-Ostrogradskij theorem:

$$\int_V \nabla \phi dV = \oint_A \phi \cdot d\underline{A}$$

The i -th component of the approximate gradient can be evaluated according to the following expression:

$$\tilde{\nabla}_i \phi = \frac{\sum_k \int_{A_k} \phi dA_i}{V_p}$$

A_i is the i -th component of the surface vector in Descartes system.

The approximate Laplacian

$$\Delta \phi = \nabla \cdot \nabla \phi$$

When calculating the discrete approximation of the operator the gradient must be interpolated onto the face centroids. This is denoted by $\langle \rangle$ in the following formula:

$$\tilde{\Delta} \phi = \tilde{\nabla} \cdot \langle \tilde{\nabla}_i \phi \rangle$$

For most field variables - excepting for the pressure field - the face normal component of the gradient vector can be calculated on a more simple way: from ϕ values stored in the centers of the adjacent cells.

In this case the discrete form of the Laplacian operator can be calculated as a linear combination of ϕ_p and the neighboring ϕ values:

$$\tilde{\Delta} \phi = A_p \phi_p + \sum A_t \phi_t$$